

11 – Waves and Instabilities

Waves

Quasi-geostrophic approximation
Potential vorticity

Waves

Quasi-geostrophic approximation

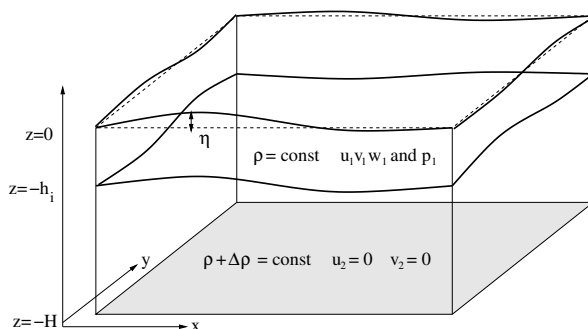
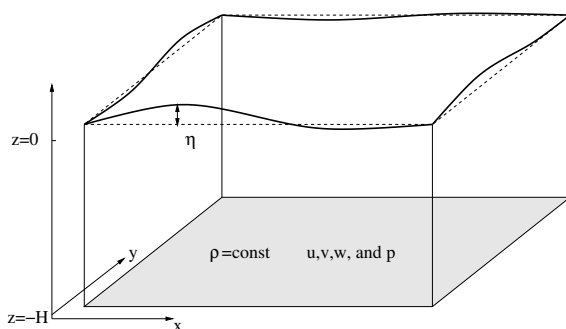
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- ▶ "barotropic model" and "baroclinic model"

$$\frac{\partial u}{\partial t} + \mathbf{u} \cdot \nabla u - fv = -g \frac{\partial h}{\partial x}, \quad \frac{\partial v}{\partial t} + \mathbf{u} \cdot \nabla v + fu = -g \frac{\partial h}{\partial y}$$

$$\frac{\partial h}{\partial t} + \frac{\partial}{\partial x}(uh) + \frac{\partial}{\partial y}(vh) = 0$$

- ▶ h is total thickness ("barotropic") or layer interface h_i ("baroclinic")
- ▶ either $g = 9.81 \text{ m/s}^2$ ("barotropic") or $g \rightarrow g\Delta\rho/\rho_0$ ("baroclinic")



- ▶ consider the layered model (first without $\mathbf{u} \cdot \nabla \mathbf{u}$ for simplicity)

$$\begin{aligned} \frac{\partial u}{\partial t} - fv &= -g \frac{\partial h}{\partial x} & , & & \frac{\partial v}{\partial t} + fu &= -g \frac{\partial h}{\partial y} \\ \frac{\partial h}{\partial t} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) &= 0 \end{aligned}$$

- ▶ take curl of momentum equation, i.e. $\partial(2.\text{eqn})/\partial x - \partial(1.\text{eqn})/\partial y$

$$\begin{aligned} \frac{\partial}{\partial x}(2.\text{eqn}) : \quad \frac{\partial}{\partial t} \frac{\partial v}{\partial x} + \frac{\partial}{\partial x}(fu) &= -g \frac{\partial}{\partial x} \frac{\partial h}{\partial y} \\ \frac{\partial}{\partial y}(1.\text{eqn}) : \quad \frac{\partial}{\partial t} \frac{\partial u}{\partial y} - \frac{\partial}{\partial y}(fv) &= -g \frac{\partial}{\partial y} \frac{\partial h}{\partial x} \end{aligned}$$

subtract both

$$\begin{aligned} \frac{\partial \zeta}{\partial t} + \frac{\partial}{\partial x}(fu) + \frac{\partial}{\partial y}(fv) &= 0 \\ \frac{\partial \zeta}{\partial t} + f \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \beta v &= 0 \end{aligned}$$

with relative vorticity $\zeta = \partial v/\partial x - \partial u/\partial y$ and with $\beta = \partial f/\partial y$

- ▶ assume small Rossby number Ro , i.e. dominant geostrophic balance

$$\begin{aligned} O(Ro) - fv &= -g \frac{\partial h}{\partial x} & , & & O(Ro) + fu &= -g \frac{\partial h}{\partial y} \\ \frac{\partial h}{\partial t} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) &= 0 \end{aligned}$$

- ▶ then $v \approx (g/f) \partial h/\partial x$ and $u \approx -(g/f) \partial h/\partial y$
- ▶ relative vorticity ζ becomes

$$\begin{aligned} \zeta &= \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \approx \frac{\partial}{\partial x} \left(\frac{g}{f} \frac{\partial h}{\partial x} \right) - \frac{\partial}{\partial y} \left(-\frac{g}{f} \frac{\partial h}{\partial y} \right) \\ &= \frac{g}{f} \left(\frac{\partial^2 h}{\partial x^2} + \frac{\partial^2 h}{\partial y^2} \right) - \frac{g}{f^2} \frac{\partial h}{\partial y} \frac{\partial f}{\partial y} \approx \frac{g}{f} \left(\frac{\partial^2 h}{\partial x^2} + \frac{\partial^2 h}{\partial y^2} \right) \end{aligned}$$

assuming $|(g\beta/f^2) \partial h/\partial y| \ll |\partial v/\partial x|$, i.e. small variations of f

$$\frac{g}{f^2} \frac{\partial h}{\partial y} \frac{\partial f}{\partial y} \sim \frac{U \Omega}{\Omega a} , \quad \frac{\partial v}{\partial x} \sim \frac{U}{L} \rightarrow \frac{U}{a} \ll \frac{U}{L} \text{ if } L \ll a$$

with Earth radius a , i.e. only length scales L smaller than a are valid

- ▶ assume small Rossby number Ro , i.e. dominant geostrophic balance

$$O(Ro) - fv = -g \frac{\partial h}{\partial x} \quad , \quad O(Ro) + fu = -g \frac{\partial h}{\partial y}$$

$$\frac{\partial h}{\partial t} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) = 0$$

- ▶ $\text{curl } \zeta \approx (g/f) (\partial^2 h / \partial x^2 + \partial^2 h / \partial y^2)$ for $Ro \ll 1$ and $L \ll a$

$$\frac{\partial \zeta}{\partial t} + f \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \beta v = 0$$

$$\frac{g}{f} \frac{\partial}{\partial t} \nabla^2 h - \frac{f}{H} \frac{\partial h}{\partial t} + \beta \frac{g}{f} \frac{\partial h}{\partial x} \approx 0$$

$$\frac{\partial}{\partial t} (\nabla^2 h - R^{-2} h) + \beta \frac{\partial h}{\partial x} \approx 0$$

with the "Rossby radius" $R = \sqrt{gH}/|f|$ and Earth radius a

- ▶ single equation in h : quasi-geostrophic potential vorticity equation valid for $Ro \ll 1$ and $L \ll a$

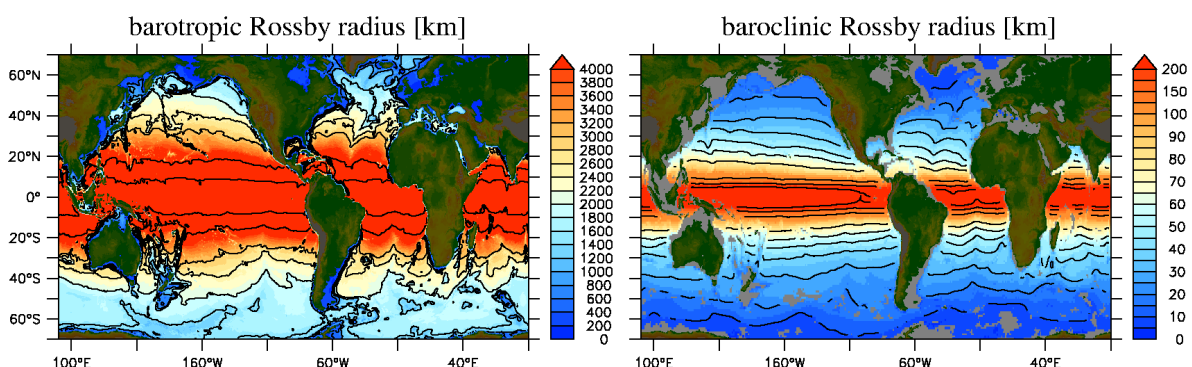
- ▶ quasi-geostrophic potential vorticity (PV) equation

$$\frac{\partial}{\partial t} (\nabla^2 h - R^{-2} h) + \beta \frac{\partial h}{\partial x} = 0$$

valid for $Ro \ll 1$ and $L \ll a$

with the "Rossby radius" $R = \sqrt{gH}/|f|$ and Earth radius a

- ▶ h is total thickness ("barotropic") or layer interface h_i ("baroclinic")
- ▶ either $g = 9.81 \text{ m/s}^2$ ("barotropic") or $g \rightarrow g\Delta\rho/\rho_0$ ("baroclinic")



- ▶ consider again the layered model, but this time with $\mathbf{u} \cdot \nabla \mathbf{u}$

$$\frac{\partial u}{\partial t} + \mathbf{u} \cdot \nabla u - fv = -g \frac{\partial h}{\partial x} \quad , \quad \frac{\partial v}{\partial t} + \mathbf{u} \cdot \nabla v + fu = -g \frac{\partial h}{\partial y}$$

- ▶ take curl of momentum equation, i.e. $\partial(2.eqn)/\partial x - \partial(1.eqn)/\partial y$

$$\begin{aligned} \frac{\partial}{\partial x}(2.eqn) : \quad & \frac{\partial}{\partial t} \frac{\partial v}{\partial x} + \frac{\partial}{\partial x} (\mathbf{u} \cdot \nabla v) + \frac{\partial}{\partial x} (fu) = -g \frac{\partial}{\partial x} \frac{\partial h}{\partial y} \\ \frac{\partial}{\partial y}(1.eqn) : \quad & \frac{\partial}{\partial t} \frac{\partial u}{\partial y} + \frac{\partial}{\partial y} (\mathbf{u} \cdot \nabla u) - \frac{\partial}{\partial y} (fv) = -g \frac{\partial}{\partial y} \frac{\partial h}{\partial x} \end{aligned}$$

subtract both

$$\begin{aligned} \frac{D\zeta}{Dt} + \frac{\partial \mathbf{u}}{\partial x} \cdot \nabla v - \frac{\partial \mathbf{u}}{\partial y} \cdot \nabla u + \frac{\partial}{\partial x} (fu) + \frac{\partial}{\partial y} (fv) &= 0 \\ \frac{D\zeta}{Dt} + \frac{\partial \mathbf{u}}{\partial x} \cdot \nabla v - \frac{\partial \mathbf{u}}{\partial y} \cdot \nabla u + f \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \beta v &= 0 \end{aligned}$$

with relative vorticity $\zeta = \partial v / \partial x - \partial u / \partial y$ and with $\beta = \partial f / \partial y$

- ▶ curl of momentum equation, i.e. $\partial(2.eqn)/\partial x - \partial(1.eqn)/\partial y$

$$\frac{D\zeta}{Dt} + \frac{\partial \mathbf{u}}{\partial x} \cdot \nabla v - \frac{\partial \mathbf{u}}{\partial y} \cdot \nabla u + f \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \beta v = 0$$

- ▶ calculate

$$\begin{aligned} \frac{\partial \mathbf{u}}{\partial x} \cdot \nabla v - \frac{\partial \mathbf{u}}{\partial y} \cdot \nabla u &= \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} + \frac{\partial v}{\partial x} \frac{\partial v}{\partial y} - \frac{\partial u}{\partial y} \frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} \frac{\partial u}{\partial y} \\ &= \frac{\partial v}{\partial x} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) - \frac{\partial u}{\partial y} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \\ &= \zeta \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) \end{aligned}$$

- ▶ now use $|\zeta| \ll |f| \rightarrow U/L \ll \Omega \rightarrow U/(\Omega L) = Ro \ll 1$

$$\frac{D\zeta}{Dt} + f \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \beta v \approx 0$$

for $Ro \ll 1$

- ▶ assume again $Ro \ll 1$, i.e. dominant geostrophic balance

$$O(Ro) - fv = -g \frac{\partial h}{\partial x} \quad , \quad O(Ro) + fu = -g \frac{\partial h}{\partial y}$$

$$\frac{Dh}{Dt} + H \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) = 0$$

- ▶ $\text{curl } \zeta \approx (g/f) (\partial^2 h / \partial x^2 + \partial^2 h / \partial y^2)$ for $Ro \ll 1$ and $L \ll a$

$$\frac{D\zeta}{Dt} + f \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) + \beta v \approx 0$$

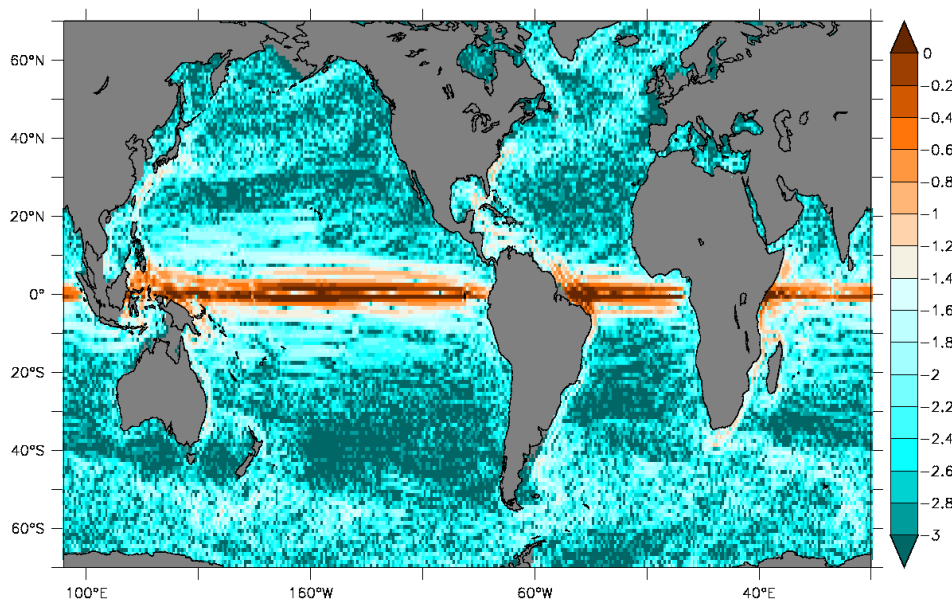
$$\frac{D}{Dt} \left(\frac{g}{f} \nabla^2 h \right) - \frac{f}{H} \frac{Dh}{Dt} + \beta \frac{g}{f} \frac{\partial h}{\partial x} \approx 0$$

$$\frac{D}{Dt} (\nabla^2 h - R^{-2} h) + \beta \frac{\partial h}{\partial x} \approx 0$$

with the "Rossby radius" $R = \sqrt{gH}/|f|$ and Earth radius a

- ▶ non-linear quasi-geostrophic PV equation: $\partial/\partial t \rightarrow D/Dt$
still valid only for $Ro \ll 1$ and $L \ll a$

- ▶ Ro as $\log_{10}(|\zeta|/|f|)$ at 100 m in high resolution ocean model



- ▶ except for equatorial ocean, $Ro = |\zeta|/|f|$ is well below 0.1

- ▶ quasi-geostrophic potential vorticity (PV) equation

$$\frac{D}{Dt} \frac{g}{f} (\nabla^2 h - R^{-2} h) + \beta v = \frac{D}{Dt} \left[\frac{g}{f_0} (\nabla^2 h - R^{-2} h) + f_0 + \beta y \right] = 0$$

with material derivative $D/Dt = \partial/\partial t + \mathbf{u} \cdot \nabla$

- ▶ "β-plane" approximation was used at $y' = y_0 + \Delta y$:

$$f(y') = f|_{y_0} + \frac{\partial f}{\partial y}|_{y_0} \Delta y + \dots \approx f_0 + \beta \Delta y \equiv f_0 + \beta y$$

with $f_0 = f|_{y_0} = \text{const}$ and $\beta = \partial f / \partial y|_{y_0} = \text{const}$

it follows that $Df/Dt = D/Dt(f_0 + \beta y) = \mathbf{u} \cdot \nabla(\beta y) = \beta v$

- ▶ quasi-geostrophic PV is approximation to full PV for single layer

$$\frac{D}{Dt} \left(\frac{\zeta + f}{h} \right) = 0$$

- ▶ full potential vorticity (PV) equation can be derived from full equations for single layer (see exercises)

- ▶ full potential vorticity equation

$$\frac{Dq}{Dt} = 0, \quad q = \frac{\zeta + f}{h}$$

- ▶ quasi-geostrophic potential vorticity equation

$$\frac{Dq}{Dt} = 0, \quad q = \frac{g}{f_0} (\nabla^2 h - R^{-2} h) + f = \zeta - \frac{f_0}{H} h + f$$

with $\zeta \approx (g/f_0) \nabla^2 h$ and $f = f_0 + \beta y$ and $R^2 = gH/f_0^2$

- ▶ approximate full q using $|\zeta| \ll |f|$ and $|\beta y| \ll |f_0|$ and $|\eta| \ll |H|$

$$\begin{aligned} \frac{\zeta + f}{h} &= \frac{\zeta}{H + \eta} + \frac{f_0}{H + \eta} + \frac{\beta y}{H + \eta} \\ &= \frac{\zeta}{H} + O(Ro) + \frac{f_0/H}{1 + \eta/H} + \frac{\beta y}{H} + O(Ro) \\ &= \frac{\zeta}{H} + \frac{f_0}{H} \left(1 - \frac{\eta}{H}\right) + \frac{\beta y}{H} + O(Ro) \\ &= (\zeta + f_0 - (f_0/H)\eta + \beta y) / H + O(Ro) \\ &= (\zeta - (f_0/H)h + f) / H + f_0/H + O(Ro) \end{aligned}$$

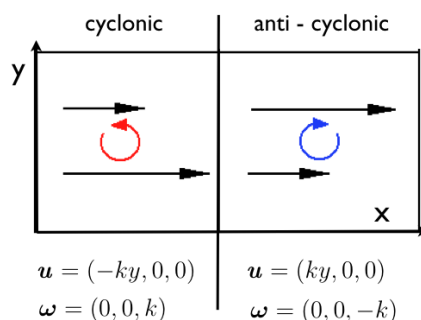
with $1/(1+x) \approx 1-x$ for small $x = \eta/H$

- ▶ potential vorticity equation for a single layer

$$\frac{Dq}{Dt} = 0 \quad , \quad q = \frac{\zeta + f}{h} \quad \text{or} \quad q = \zeta - \frac{f_0}{H}h + f$$

q is conserved for fluid parcels in single layer

- ▶ $h = \text{const}$, ζ initially zero, parcel moves northward
 f increases but $q = (f + \zeta)/h$ has to stay constant
 $\rightarrow \zeta = \partial v/\partial x - \partial u/\partial y$ decreases \rightarrow anticyclonic rotation
- ▶ $h = \text{const}$, ζ initially zero, parcel moves southward
 $\rightarrow \zeta = \partial v/\partial x - \partial u/\partial y$ increases \rightarrow more cyclonic rotation



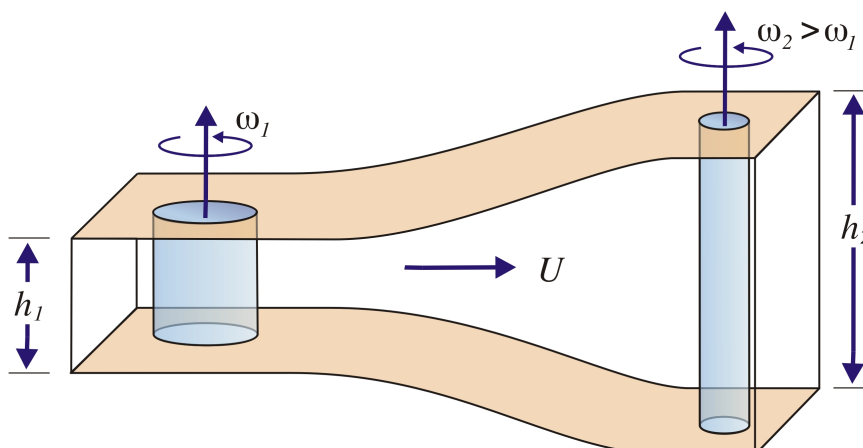
$u = -ay$, $v = 0 \rightarrow \zeta = a > 0$: cyclonic (anticlockwise) rotation
 $u = +ay$, $v = 0 \rightarrow \zeta = a < 0$: anticyclonic (clockwise) rotation

- ▶ potential vorticity equation for a single layer

$$\frac{Dq}{Dt} = 0 \quad , \quad q = \frac{\zeta + f}{h} \quad \text{or} \quad q = \zeta - \frac{f_0}{H}h + f$$

q is conserved for fluid parcels in single layer

- ▶ $f = \text{const}$, ζ initially zero, parcel moves to deeper water
 $\rightarrow \zeta = \partial v/\partial x - \partial u/\partial y$ increases \rightarrow cyclonic rotation

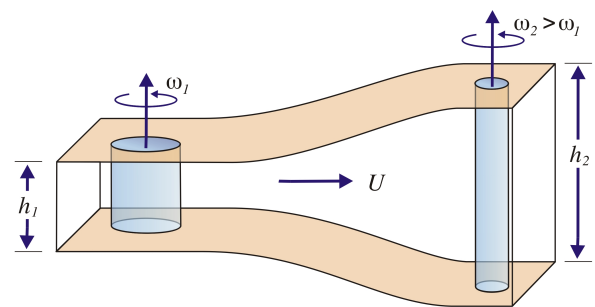


- ▶ quasi-geostrophic potential vorticity equation

$$\frac{Dq}{Dt} = 0 \quad , \quad q = \frac{g}{f} (\nabla^2 h - R^{-2}h) + f_0 + \beta y$$

q is (approximately) conserved in single layer for $Ro \ll 1$

- ▶ $\zeta = (g/f)\nabla^2 h$ is relative vorticity
- ▶ $-(g/f)R^{-2}h$ is stretching vorticity
- ▶ $f = f_0 + \beta y$ is planetary vorticity
- ▶ h is streamfunction for the quasi-geostrophic flow



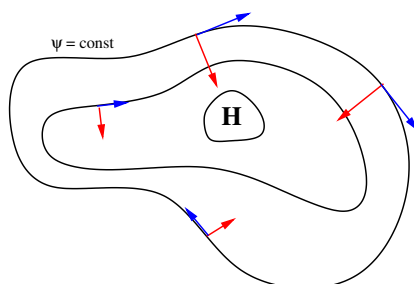
- ▶ quasi-geostrophic potential vorticity equation

$$\frac{Dq}{Dt} = 0 \quad , \quad q = \frac{g}{f} (\nabla^2 h - R^{-2}h) + f_0 + \beta y$$

q is (approximately) conserved in single layer for $Ro \ll 1$

- ▶ $\psi = gh/f_0$ is streamfunction for the quasi-geostrophic flow

$$u \approx -\frac{g}{f_0} \frac{\partial h}{\partial y} = -\frac{\partial \psi}{\partial y} \quad , \quad v \approx \frac{g}{f_0} \frac{\partial h}{\partial x} = \frac{\partial \psi}{\partial x}$$



$$\begin{aligned} \mathbf{u} &= \begin{pmatrix} -\partial\psi/\partial y \\ \partial\psi/\partial x \end{pmatrix} \\ &= \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \times \begin{pmatrix} \partial\psi/\partial x \\ \partial\psi/\partial y \\ 0 \end{pmatrix} = \mathbf{k} \times \nabla\psi \end{aligned}$$

- ▶ \mathbf{u} (blue arrow): anti-clockwise rotation of $\nabla\psi$ (red arrow) by 90°